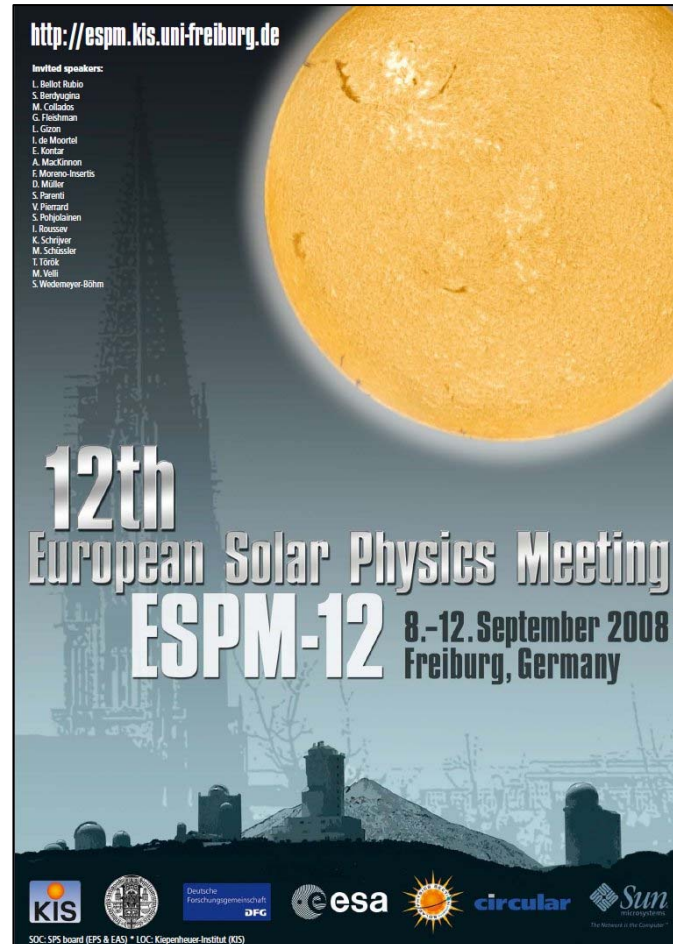


12th European Solar Physics Meeting

8 - 12 September 2008

Freiburg, Germany



Electronic proceedings

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Tuesday 16:30-17:00

Particle Acceleration and Energy Release in RHESSI Era

Kontar, E.

University of Glasgow

Since high energy emission (X-rays and gamma-rays) represents optically-thin radiation from energetic particles, it is a relatively straightforward, and hence extremely valuable, tool in the diagnostic study of flare-accelerated electrons and ions at the Sun. The observed X-ray/gamma-ray flux is fundamentally a convolution of the cross-section for the emission process(es) in question with the distribution function(s) of accelerated particles, which are in turn functions of energy, direction, spatial location and time. To address the key problems of particle acceleration, propagation as well as energy release one needs to infer as much information as possible on the particle distribution function, through a de-convolution of this fundamental relationship.

This review presents recent observational progress toward the understanding of energy release and particle acceleration using spectroscopic, imaging and polarization measurements, primarily from the Ramaty High Energy Solar Spectroscopic Imager (RHESSI). Previous conclusions regarding the energy, angular (pitch angle) and spatial distributions of energetic electrons and ions in solar flares are critically reviewed. The diagnostics of radiation processes, particle transport, and acceleration, using both spectroscopic and imaging techniques will be discussed. The unprecedented quality of the RHESSI data in combination with novel data analysis techniques have revealed previously unknown details of energetic particle distributions and imposed new challenging constraints on the particle acceleration.



University
of Glasgow

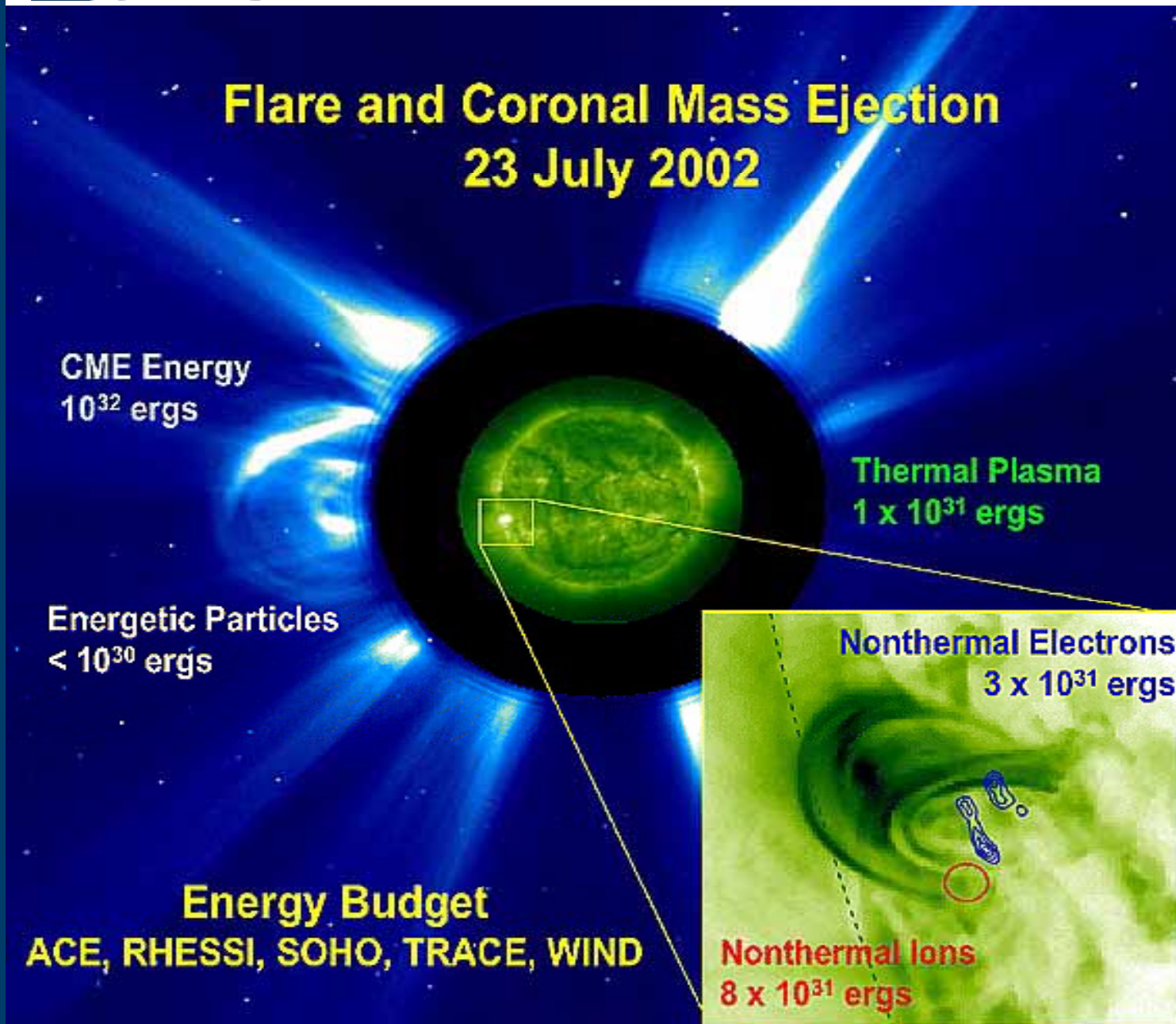
PARTICLE ACCELERATION AND ENERGY RELEASE IN RHESSI ERA

Eduard Kontar

*Department of Physics and Astronomy
University of Glasgow, UK*

12th European Solar Physics Meeting, Freiburg, Germany

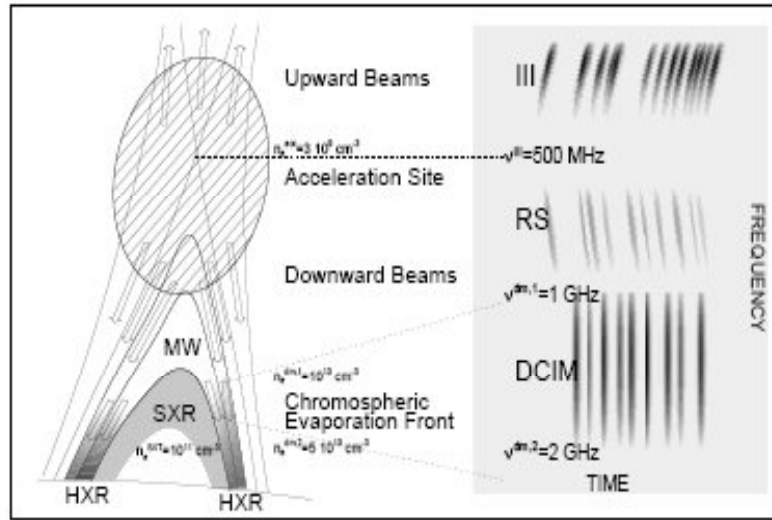
From Emslie
et al., 2004



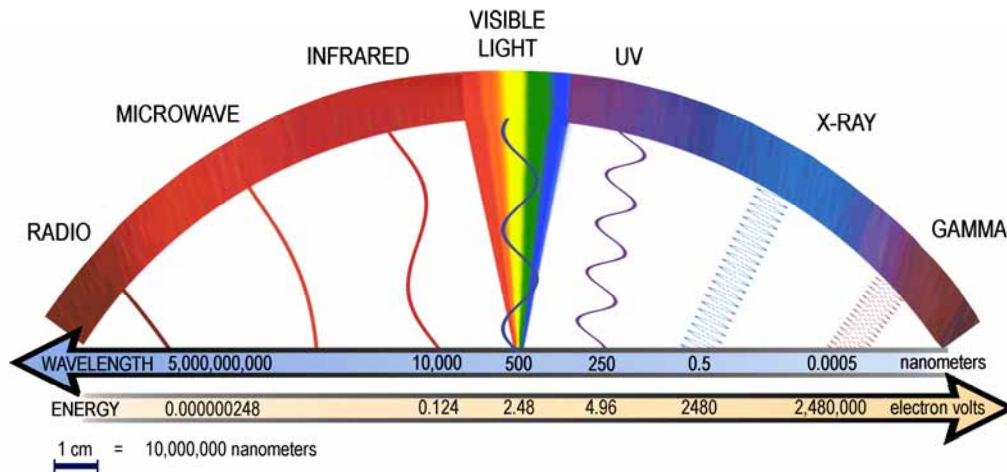
Free magnetic
energy
 $\sim 2 \times 10^{32}$ ergs



Observations of energetic particles



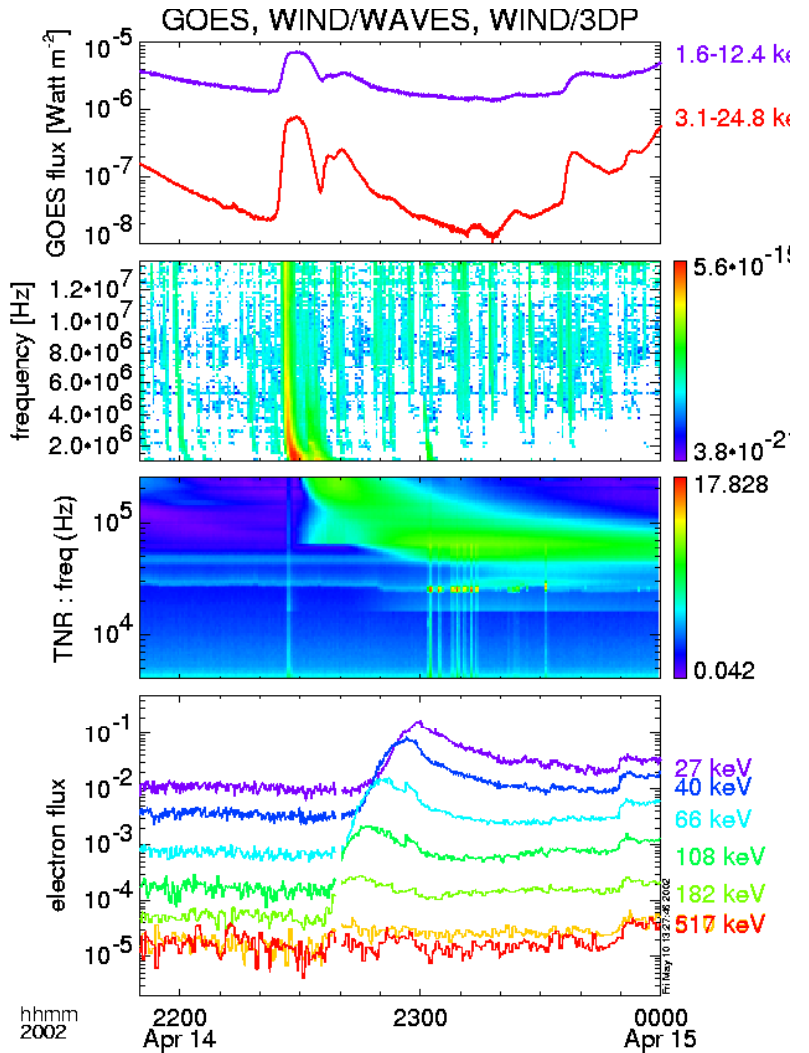
Aschwanden and Benz, 1997



X-rays

radio waves

Particles 1 AU



Krucker et al, 2001

Solar corona $T \sim 10^6 \text{ K} \Rightarrow 0.1 \text{ keV per particle}$

Flaring region $T \sim 4 \times 10^7 \text{ K} \Rightarrow 3 \text{ keV per particle}$

Flare volume $10^{27} \text{ cm}^3 \Rightarrow (10^4 \text{ km})^3$

Plasma density 10^{10} cm^{-3}

Photons up to $> 100 \text{ MeV}$

Electron energies $> 10 \text{ MeV}$

Proton energies $> 100 \text{ MeV}$

Number of energetic electrons 10^{36} per second

Typical timescale $< 0.1 \text{ sec}$

Particle spectrum is a power-law (or combination of a few)

Anti-correlation between spectral index and particle flux

Enhanced abundances of ions

Acceleration of ions/electron in different locations?

Large solar flare releases about 10^{29} ergs/s

(about half energy in energetic electrons)

1 megaton of TNT is equal to about 4×10^{22} ergs.

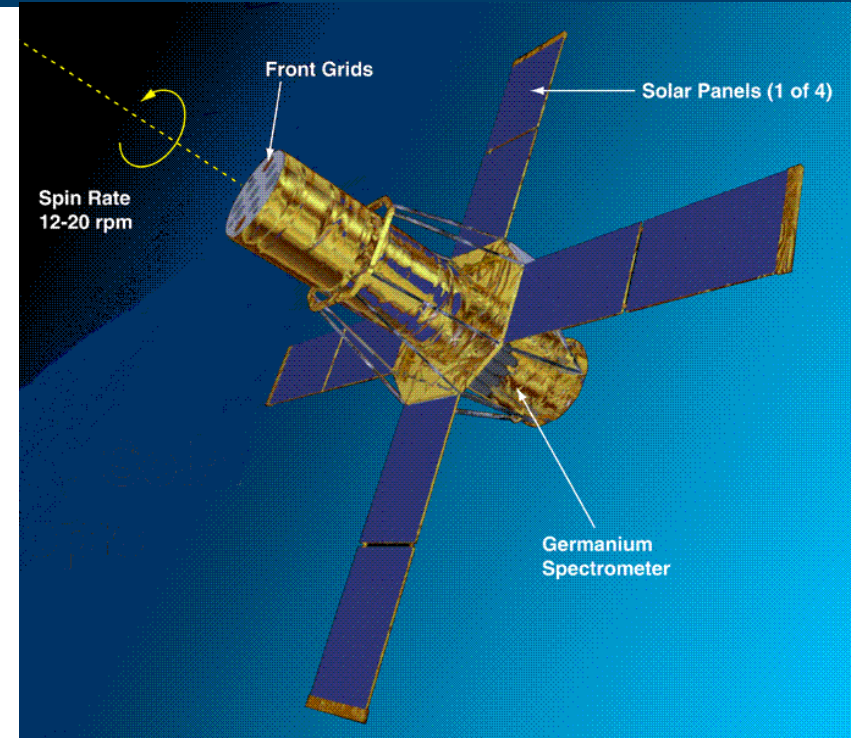
Ramaty High Energy Solar Spectroscopic Imager

is a NASA-led mission launched in February 2002

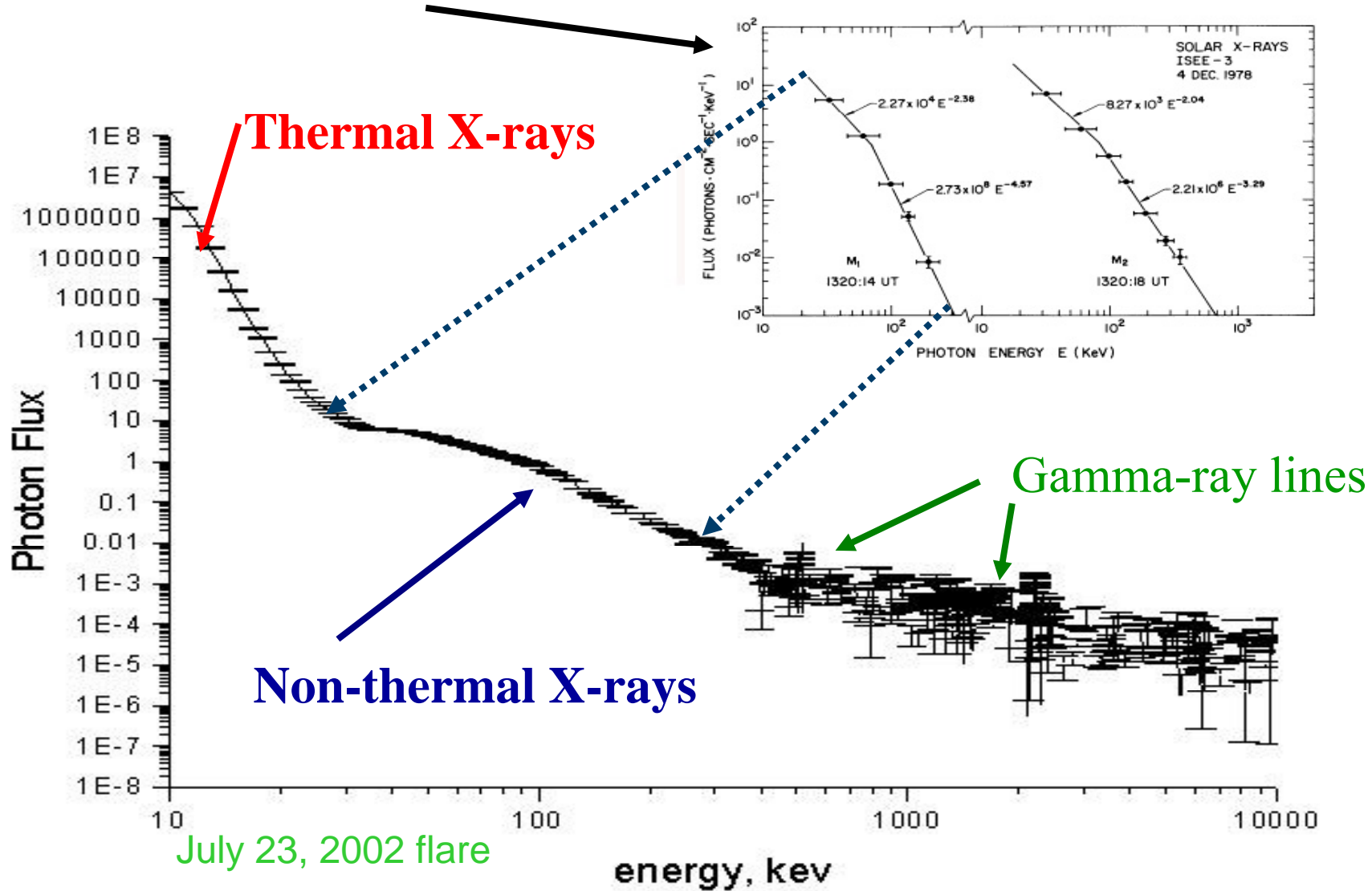
RHESSI is designed to investigate particle acceleration and energy release in solar flares through imaging and spectroscopy of hard X-ray and gamma-rays in the range from 3 keV up to 17 MeV (Lin et al 2002).

Spectroscopy: 9 Ge detectors with energy resolution around 1 keV;

Imaging: rotating modulating collimators allowing angular resolution down to 2.3 arcsec; **Imaging spectroscopy**



pre-RHESSI X-ray spectra (Kane et al, 1982)



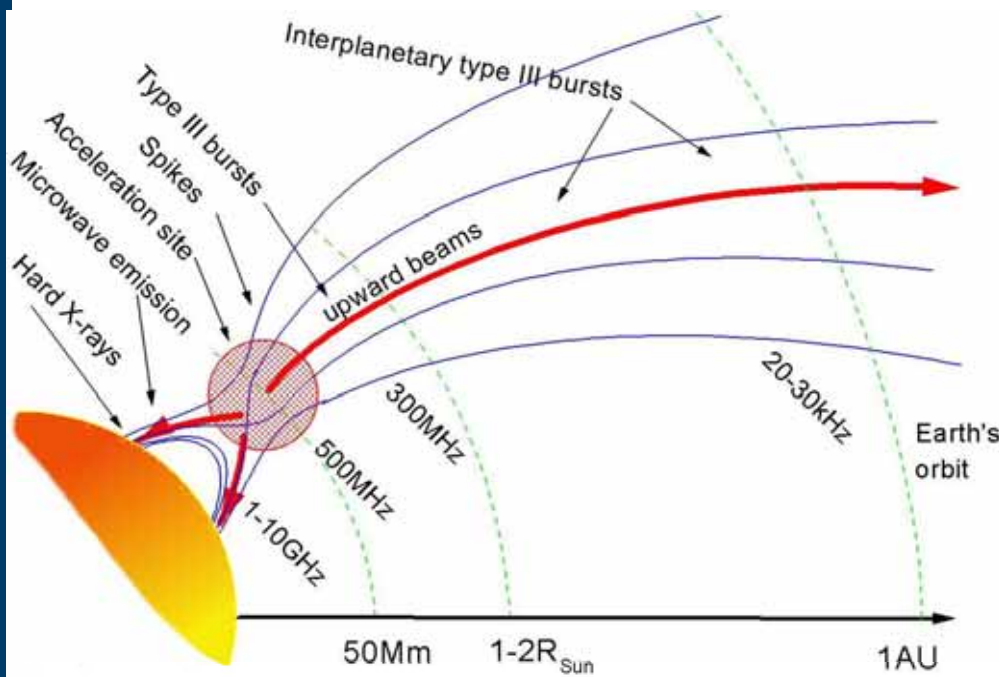
Ramaty High Energy Solar Spectroscopic Imager (RHESSI) spectrum

Observed X-rays

Unknown electron distribution

Emission cross-sections

$$I(\epsilon, \Omega, t) = \int_{\ell} \int_{\Omega'} \int_{\epsilon}^{\infty} n(\mathbf{r}) \bar{F}(E, \Omega', \mathbf{r}, t) Q(\Omega, \Omega', \epsilon, E) dE d\Omega' d\ell,$$



$F(E, \Omega, r) = ???$

- 1) What is the **energy** distribution, $F(E)$?
- 2) What is the **angular** distribution, $F(E, \Omega)$?
- 3) What is **spatial** distribution, $F(E, r)$?

For spatially integrated spectrum:
$$I(\epsilon) = \frac{1}{4\pi R^2} \bar{n}V \int_{\epsilon}^{\infty} \bar{F}(E) Q(\epsilon, E) dE,$$

Thin-target case: For the electron spectrum $\mathbf{F(E) \sim E^{-\delta}}$,

a) Electron-ion bremsstrahlung (free-free emission)

Dominant process for energies $\sim 10 - 400$ keV

the photon spectrum is $\mathbf{I(\epsilon) \sim \epsilon^{-\delta-1}}$

In the simplest form Kramers' approximation:
$$Q(\epsilon, E) = Z^2 \frac{\sigma_0}{\epsilon E},$$

b) Electron-electron bremsstrahlung (free-free emission)

Dominant process for energies above 400 keV

the photon spectrum is $\mathbf{I(\epsilon) \sim \epsilon^{-\delta}}$

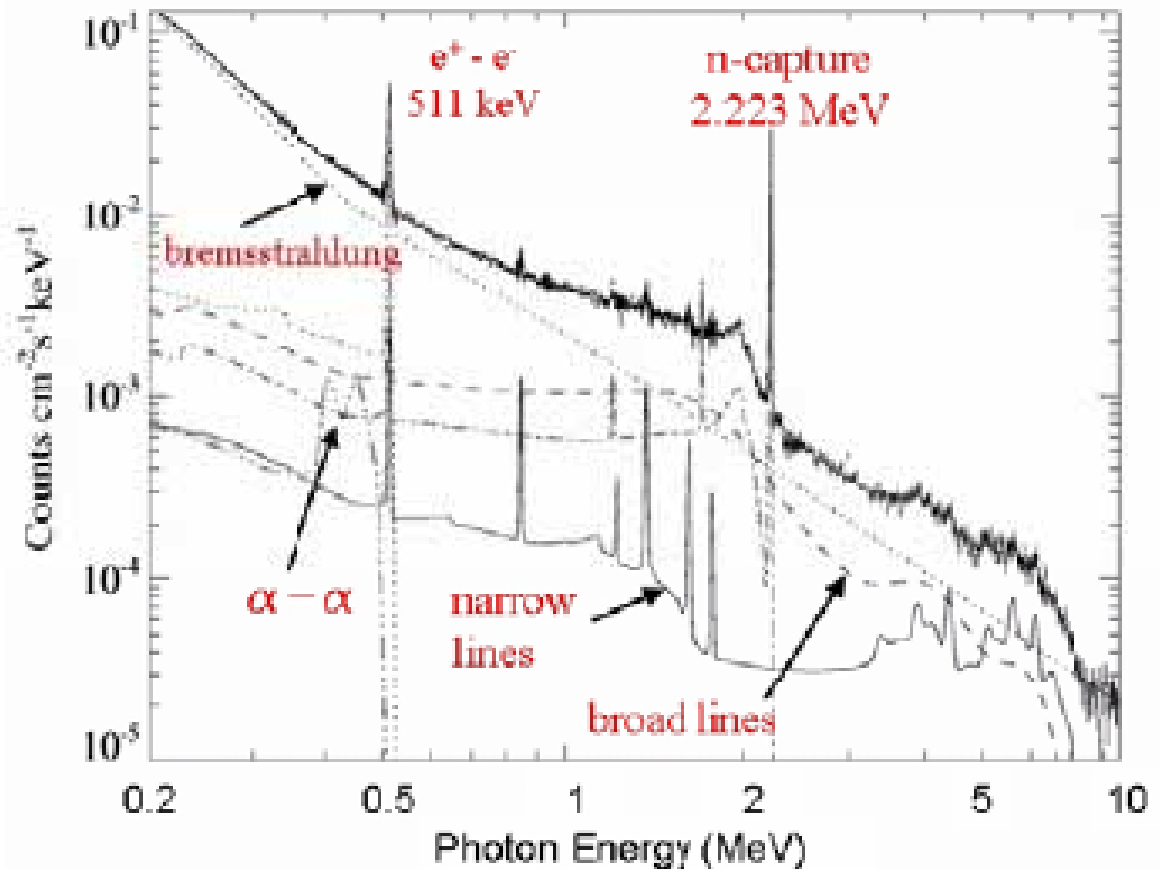
c) Recombination emission (free-bound emission)

Could be dominant process for energies up to 20 keV

the photon spectrum is **shifted by ionisation potential** and $\mathbf{I(\epsilon) \sim \epsilon^{-\delta-2}}$

(The process requires high temperatures and detailed ionisation calculations)

From Murphy and Share, 2004

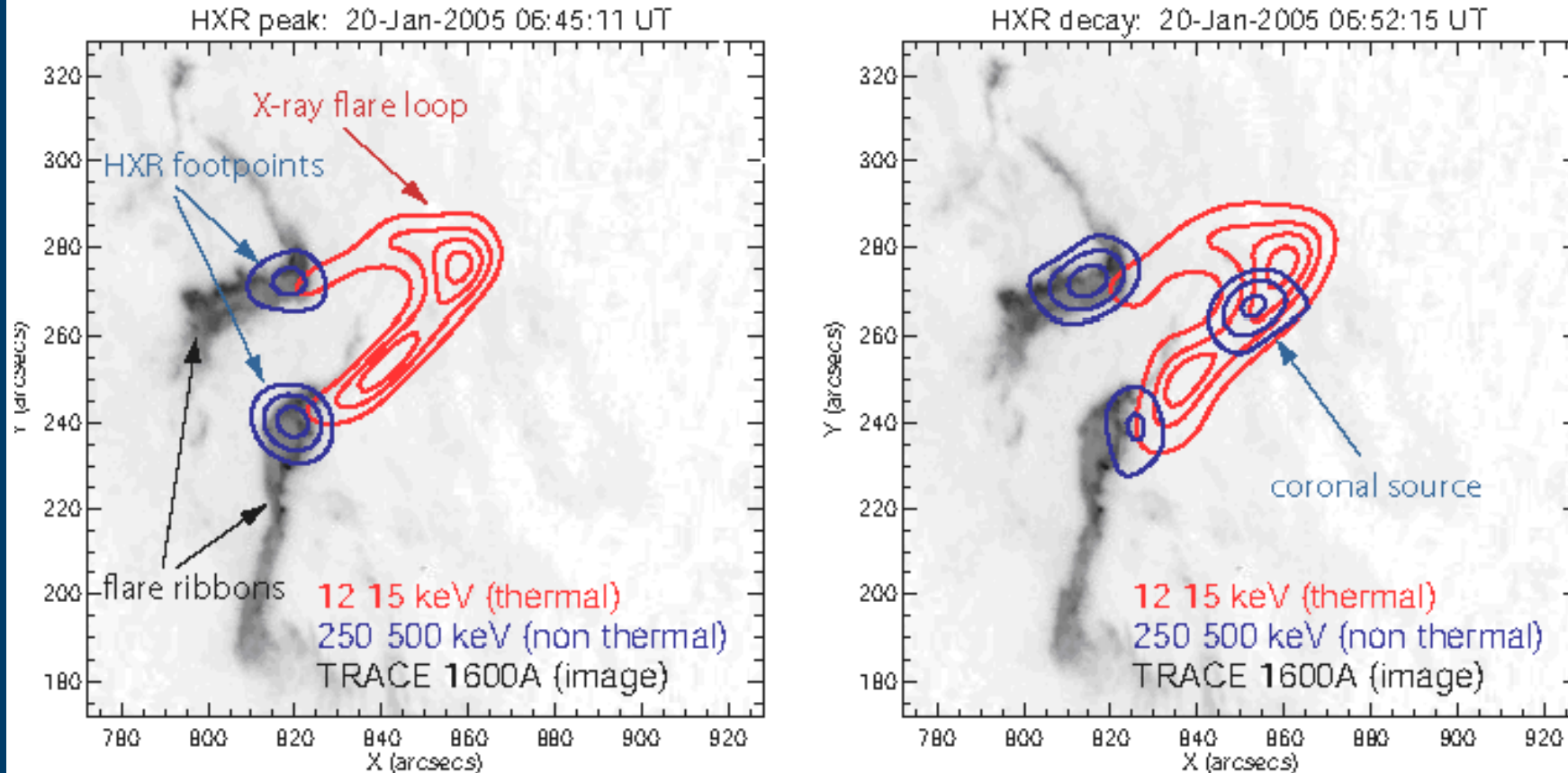


a) narrow-gamma lines

Accelerated protons and alpha particles

b) broad-gamma lines and gamma-ray-continuum

Accelerated heavy ions and unresolved lines and Compton scattering

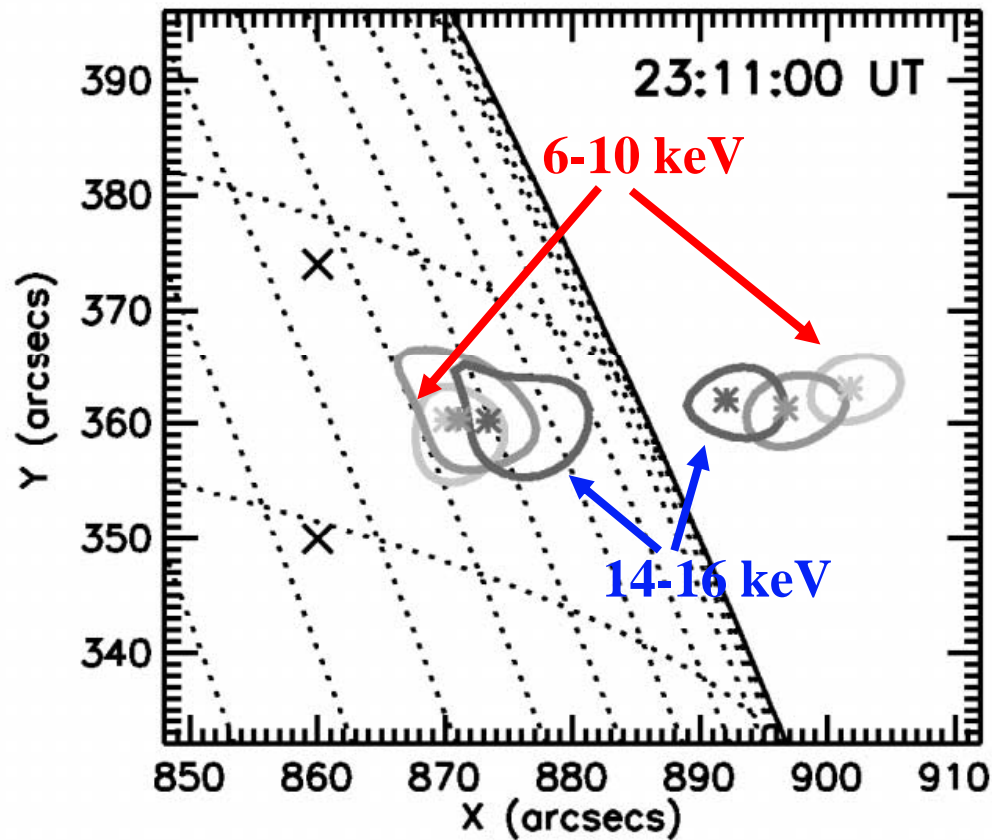


Krucker et al, 2007

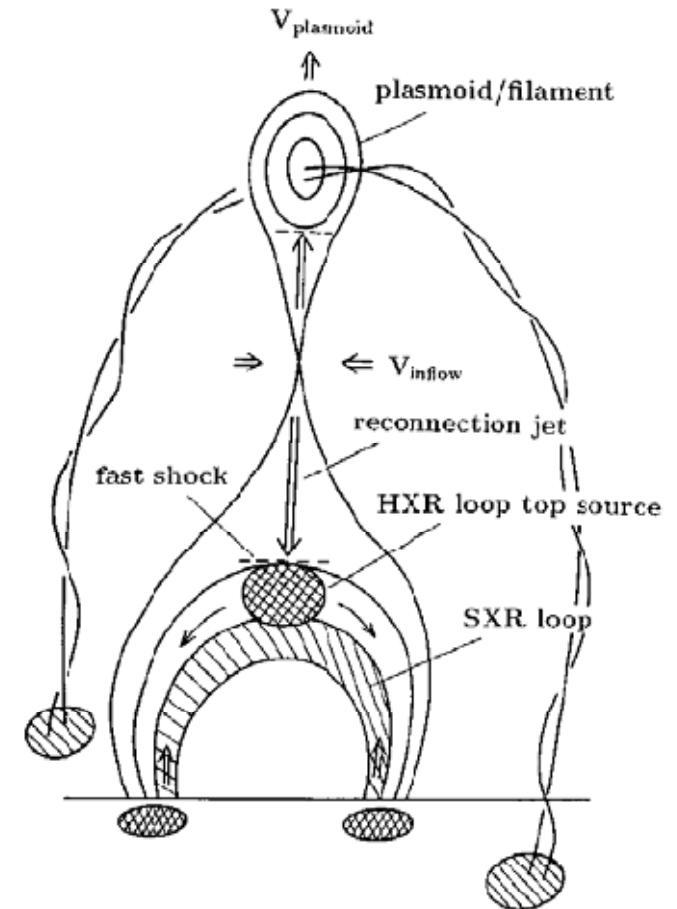
Solar flare geometry in X-rays:

Soft loop-top source and hard X-ray footpoints (Krucker & Lin 2002; Emslie et al, 2003)

What is the origin of nonthermal coronal X-ray sources?

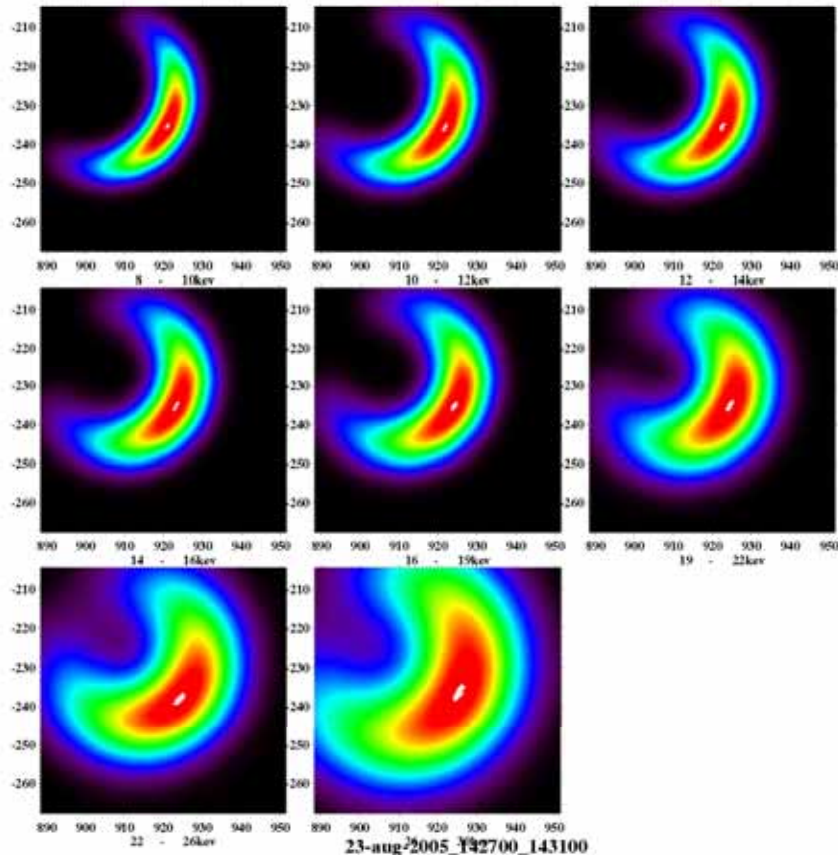


Liu et al, 2004



Shibata, 1996

Do we observe magnetic reconnection?



Number of particles in acceleration region

$$N = n A (2L)$$

Hard X-ray intensity $I(\varepsilon)$ is proportional to the rate of injection of electrons at energy $E_0 > \varepsilon$:

$$dN/dt (> \varepsilon) \sim 10^{34} I(\varepsilon)$$

Specific acceleration rate

$$\gamma = (1/N) dN/dt \text{ (particles } s^{-1} \text{ per particle)}$$

From Xu et al, 2007, Emslie et al, 2008

$$E_D = \frac{e^3 n_e \ln \Lambda}{6\pi\epsilon_0^2 kT}$$

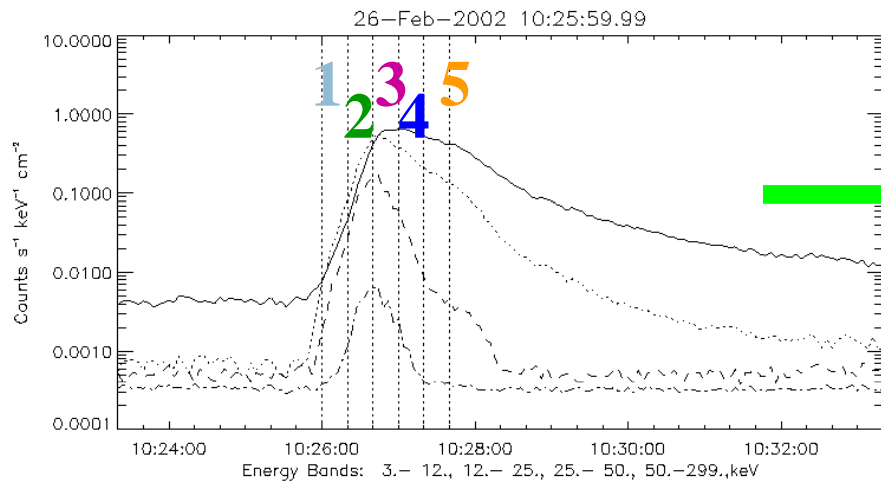
Dreicer field
(Dreicer, 1959)

The size and density of the acceleration region, plus the hard X-ray brightness, can be used to determine the **specific acceleration rate** (particles s⁻¹ particle⁻¹) – values are $\sim (0.1 - 5) \times 10^{-3}$

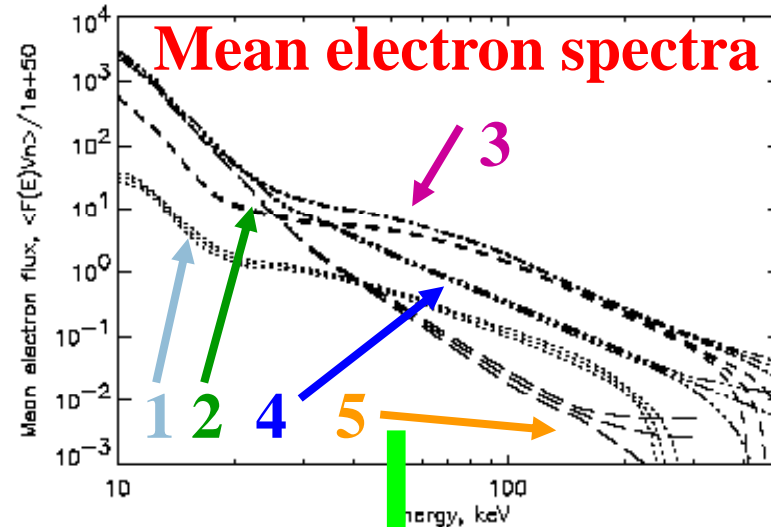
Consistency with **sub-Dreicer models** (e.g. Kuijpers (1981), Heyvaerts (1981), Holman (1985), etc) require a very narrow range of accelerating electric fields

For **super-Dreicer** current sheet acceleration (e.g. Martens (1998), Litvinenko (1996, 2003), Fletcher & Petkaki, 1997, Mori et al, 1998, Browning & Dalla, 2007), the specific acceleration rate is determined by the aspect ratio of the current sheet.

For **stochastic acceleration models** (e.g. Miller 1991 etc), values for the specific acceleration rate are generally consistent with the data, but more simulations are needed.

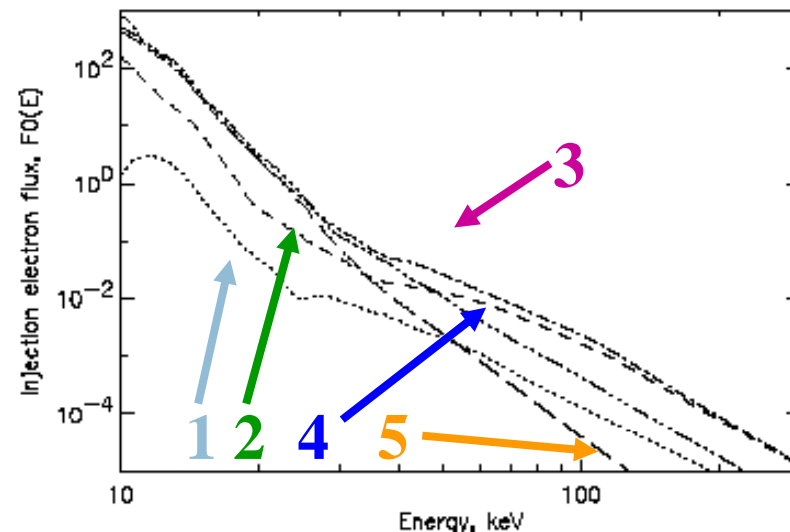


(Model independent)



(Model dependent)

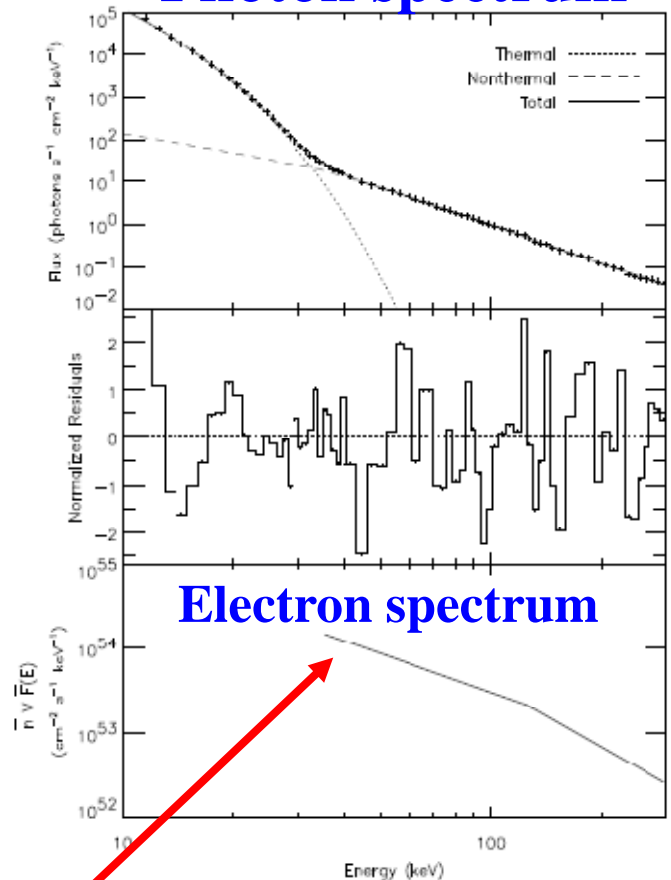
Accelerated (injected) spectra



- Deviations from power-laws; Spectral features inconsistent with simple models
- High/low energy cutoffs in the electron spectra
- Possibility to study acceleration and propagation effects on non-uniform plasma ionisations, return currents, etc

- Requiring that the assumed thermal emission dominates over non-thermal emissions, Sui *et al.* (2005) and that a **low energy cutoff of > 24 keV should be present.**
- Assuming "theoretical Neupert effect" to be satisfied, Veronig *et al.* (2005) conclude that the low-energy **cutoffs should be between 10 keV and 30 keV for four flares** analysed in their paper.
- Hannah *et al.* (2008) have used an empirical relation between the observed parameters of the photon power-law and the low-energy cutoff of the electron distribution, and have found that **the low-energy cutoffs in 25000 microflare events could range from 9 to 16 keV** with the median being around 12 keV.

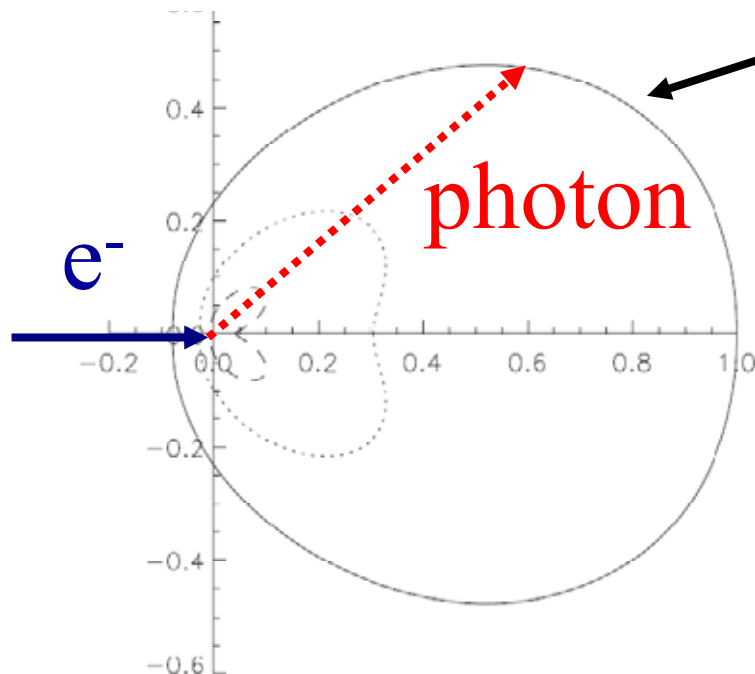
Photon spectrum



Low-energy cutoff in nonthermal electron spectrum

The observed photon flux spectrum at the Earth:

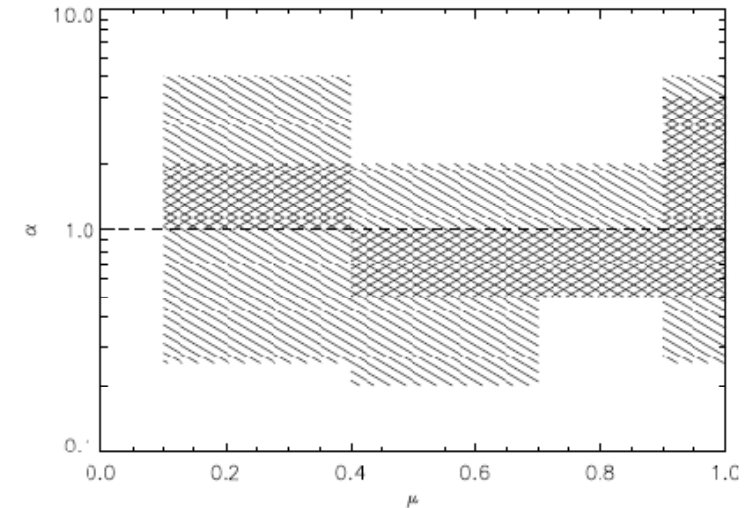
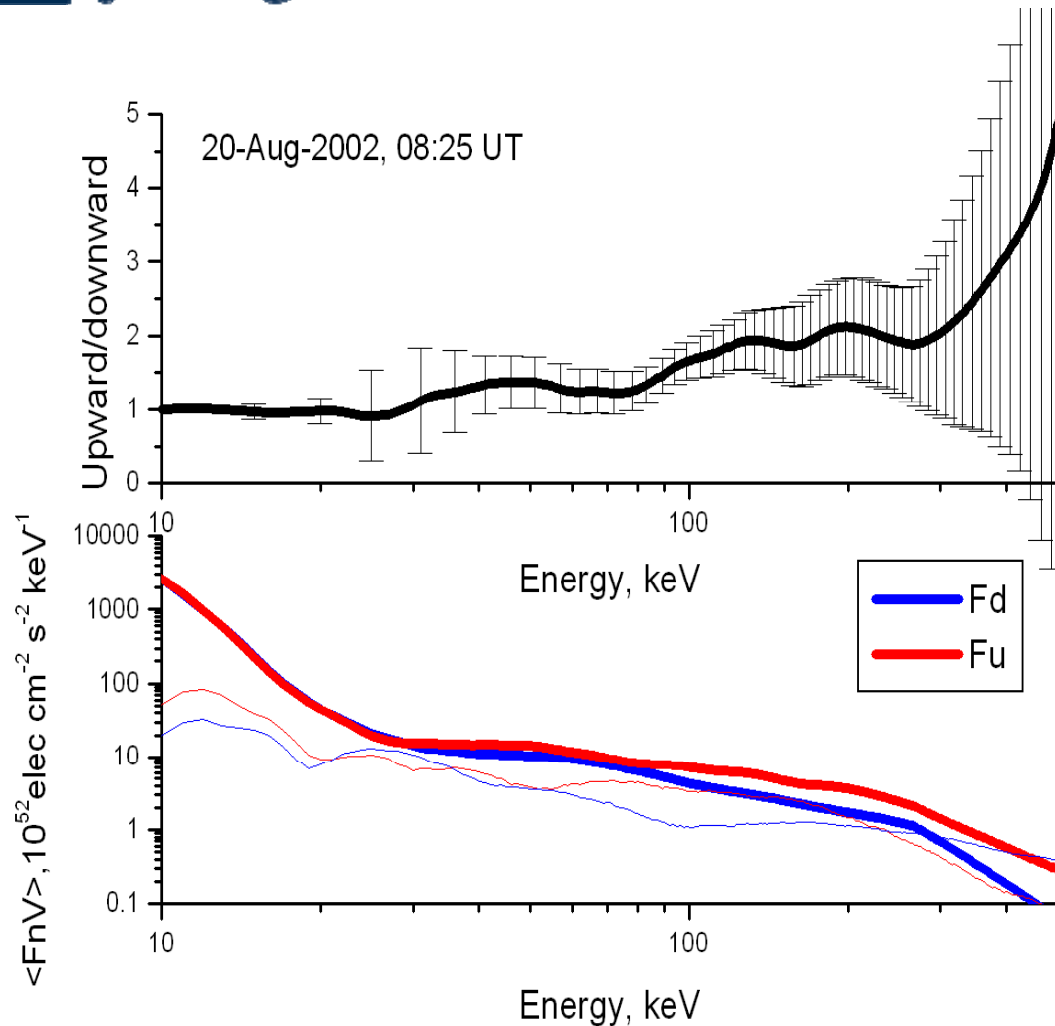
$$I(\epsilon) = \frac{\bar{n}V}{4\pi R^2} \int_{\Omega'} \int_{\epsilon}^{\infty} \bar{F}(E, \Omega') Q(\Omega, \Omega', \epsilon, E) dE d\Omega',$$



Because we observe 1D photon spectrum 3D character of electron distribution is often ignored and $F(E, \Omega)$ is assumed **isotropic**

How to measure electron anisotropy?

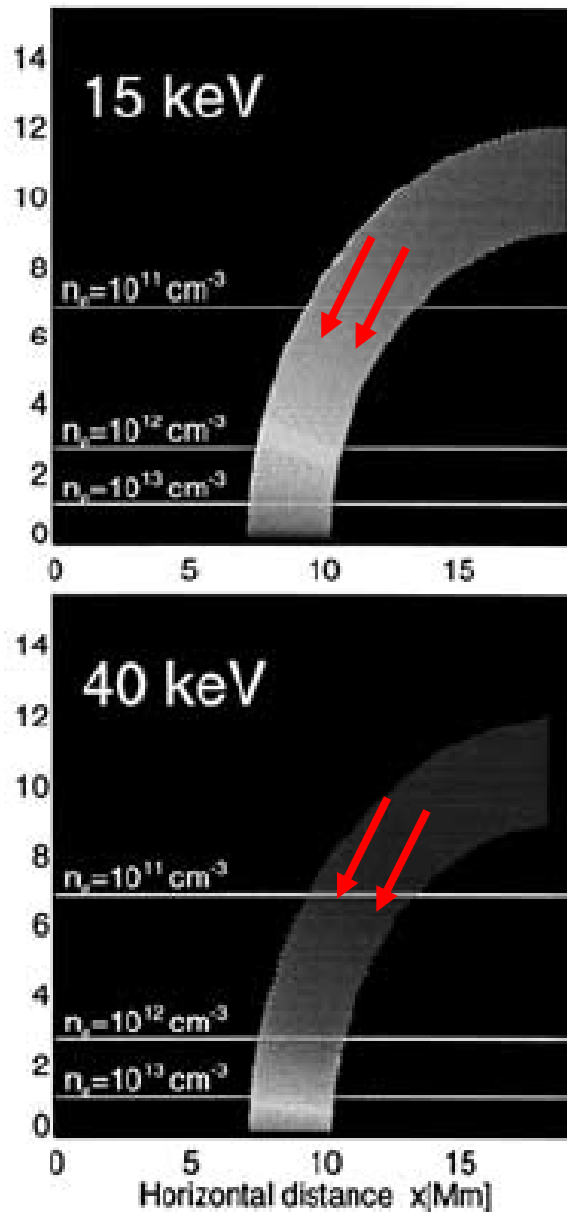
- 1) Stereoscopic X-ray observations (*Kane et al, 1982 etc*)
- 2) X-ray polarization (e.g. *McConnell et al (2003);*)
- 3) Centre-to-limb variations in solar flares (*Ohki (1969), Pinter (1969) at 10 keV, Datlowe et al. (1977) etc*)
- 4) Albedo as a probe of electron angular distribution (*Kontar & Brown 2006; Kasparova et al 2007*)



Consistent with isotropic distribution (e.g. Kane et al, 1988; Kontar&Brown, 2006, Kasparova et al, 2007)

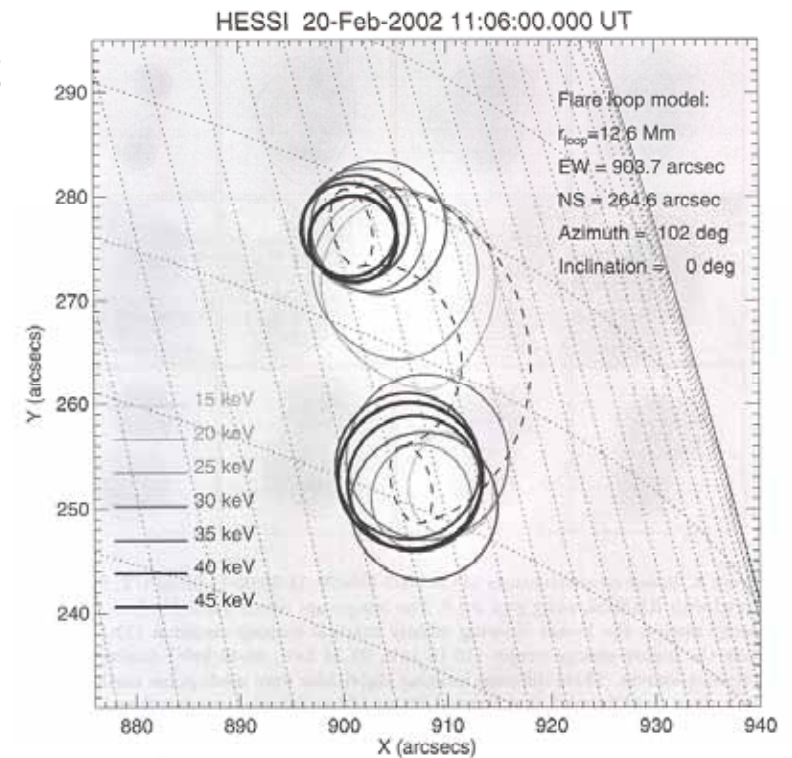
Collisional scattering and return current effects cannot explain the isotropy of electron distribution

=> The angular distribution found is **inconsistent with downward beamed distributions**



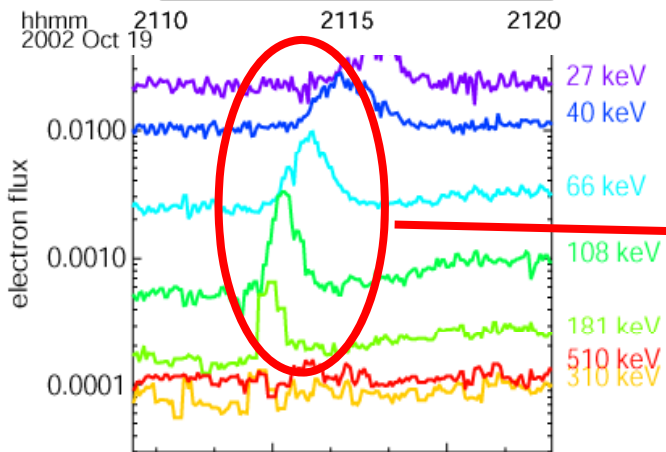
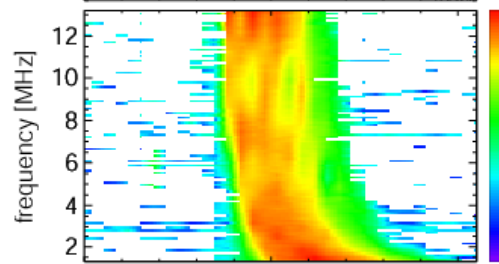
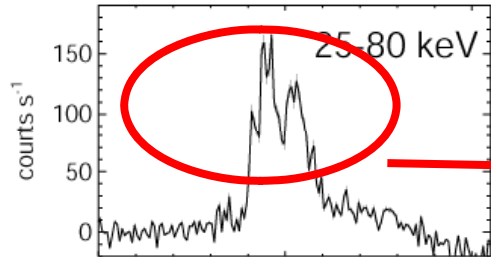
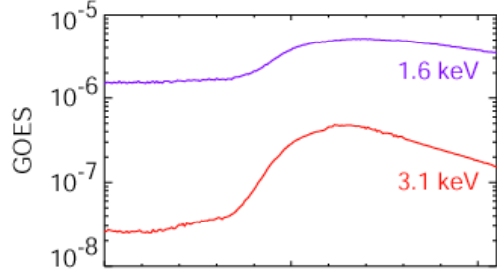
Aschwanden et al, 2002

Higher energy sources appear lower in the chromosphere (consistent with simple collisional transport)
=> downward electron beaming



Timing analysis (e.g. Aschwanden et al, 1995) also suggests the beaming of electrons

= > electrons beam downwards as in a classical scenario?

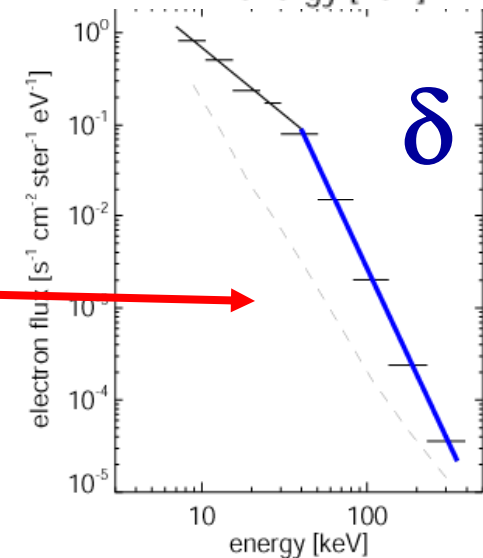
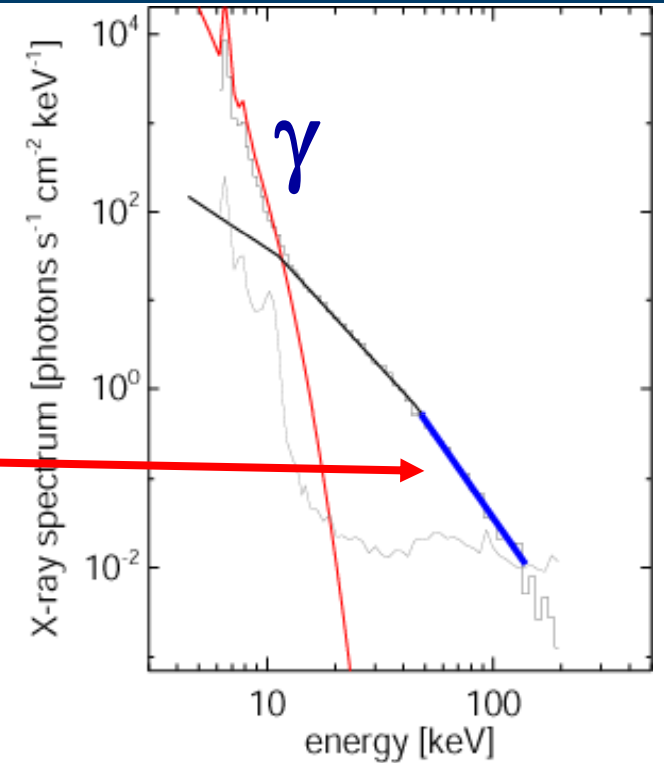


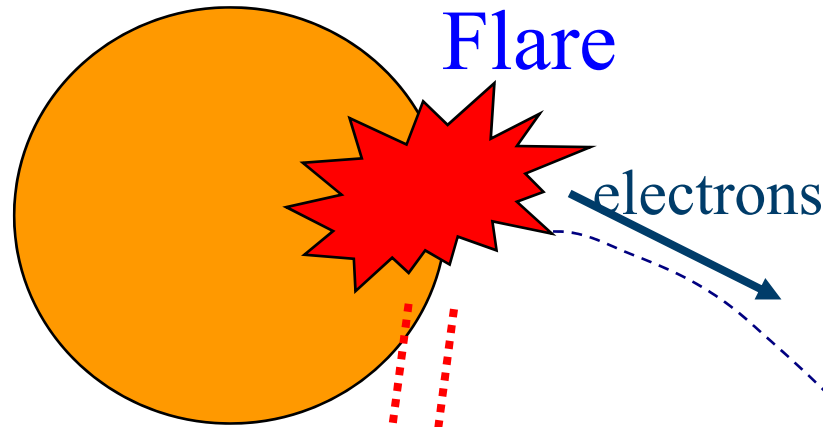
hhmm
2002 Oct 19

X-ray spectra from RHESSI

Electron spectra at 1AU from WIND

From Krucker et al 2007

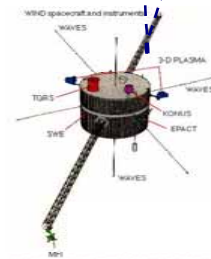




From the analysis of 16 “scatter-free” events (Lin, 1985; Krucker et al, 2007) :
 Although there is correlation between the total number of electrons at the Sun (thick-target model estimate) the spectral indices do not match either **thick-target or thin-target models**.

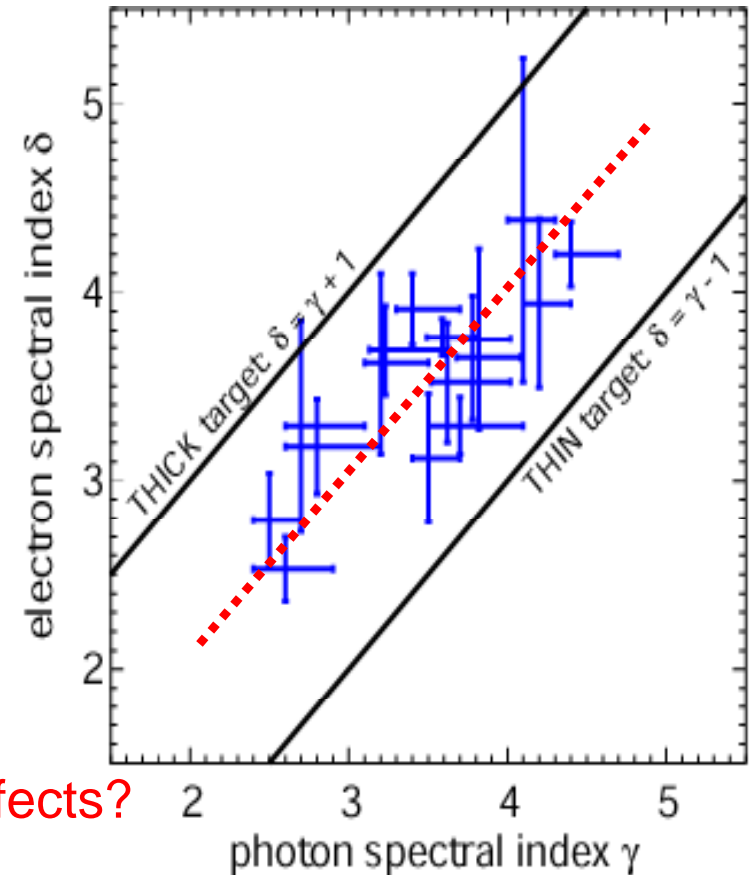


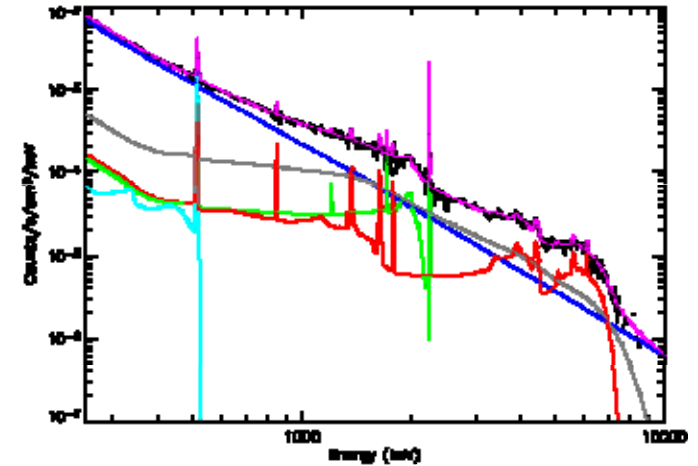
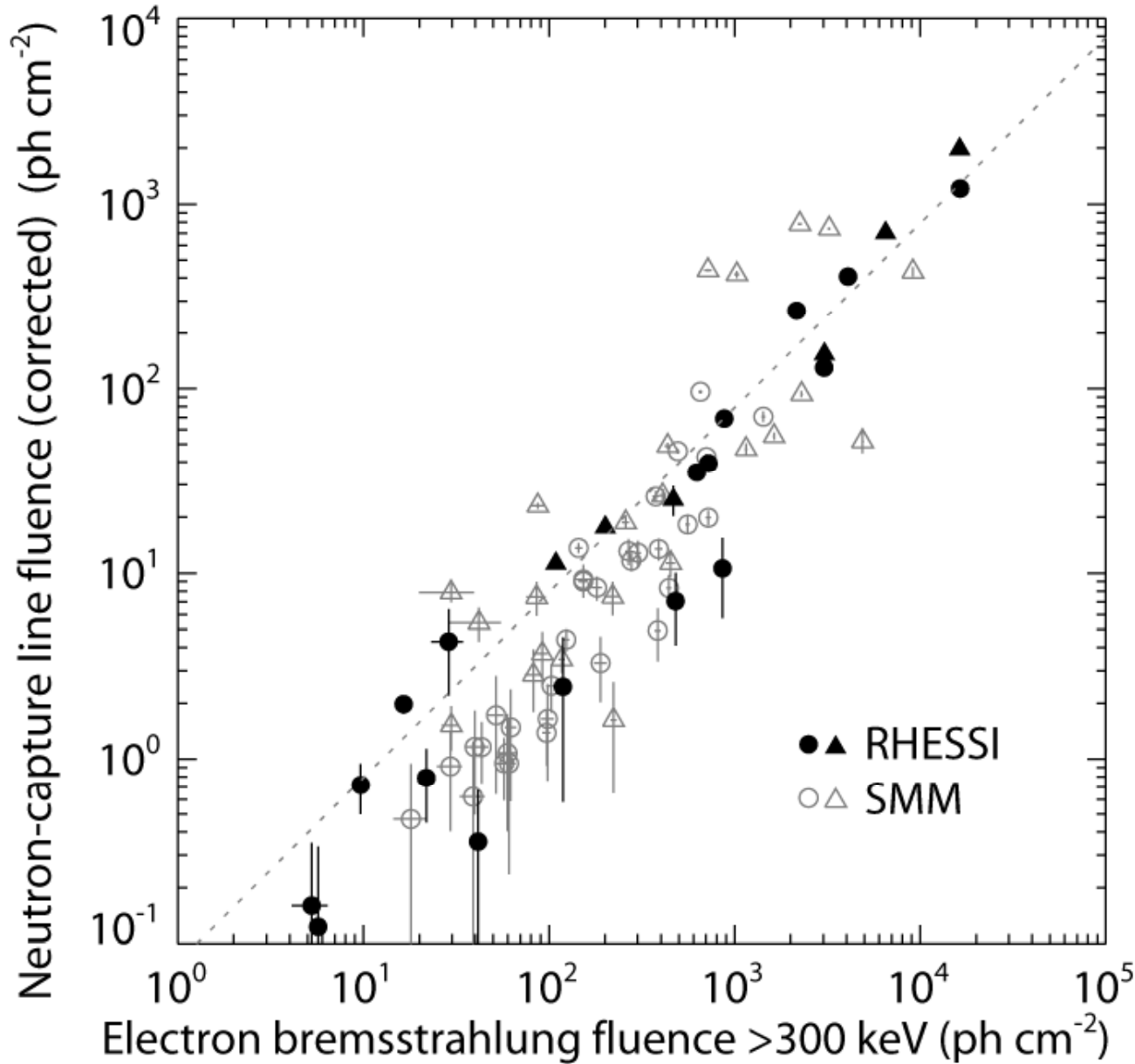
RHESSI



WIND

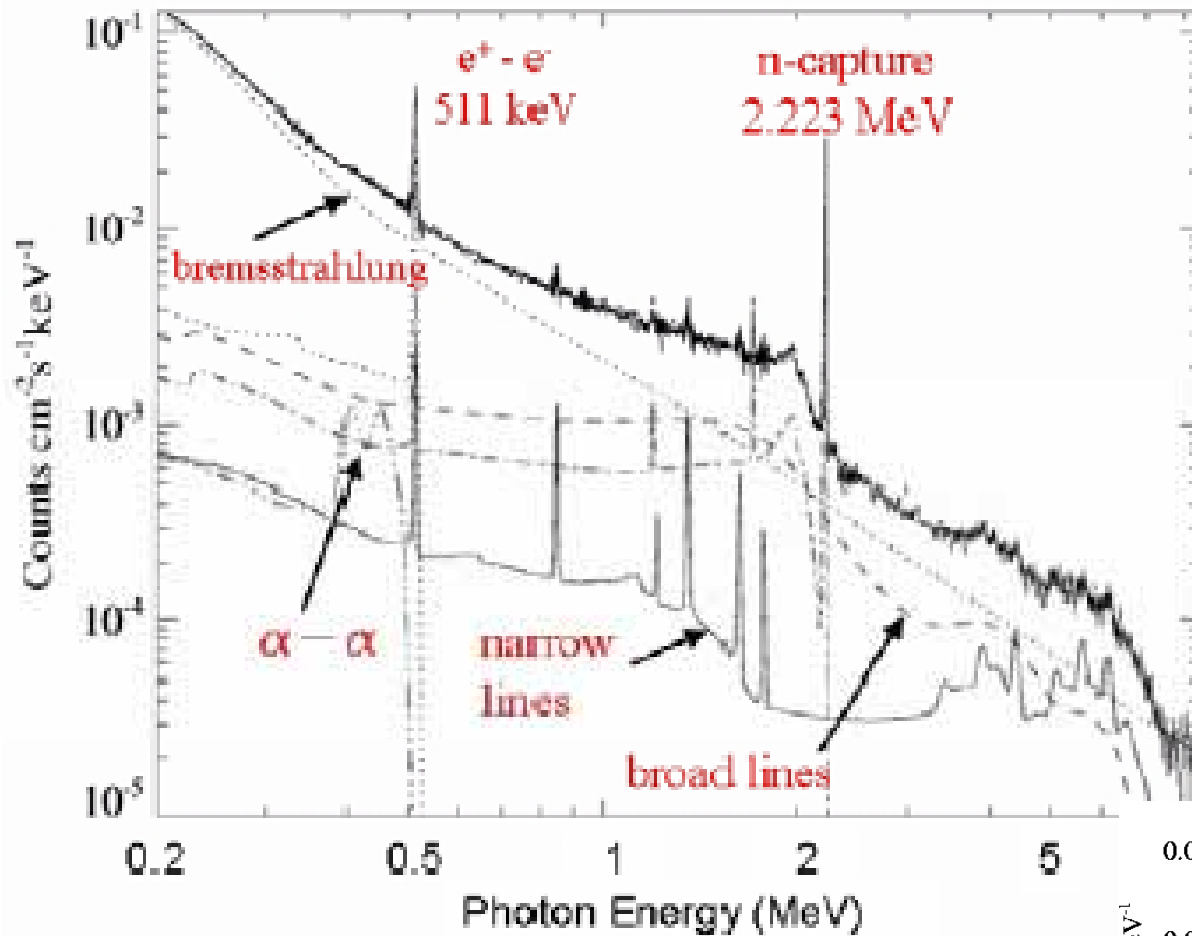
Acceleration or transport effects?



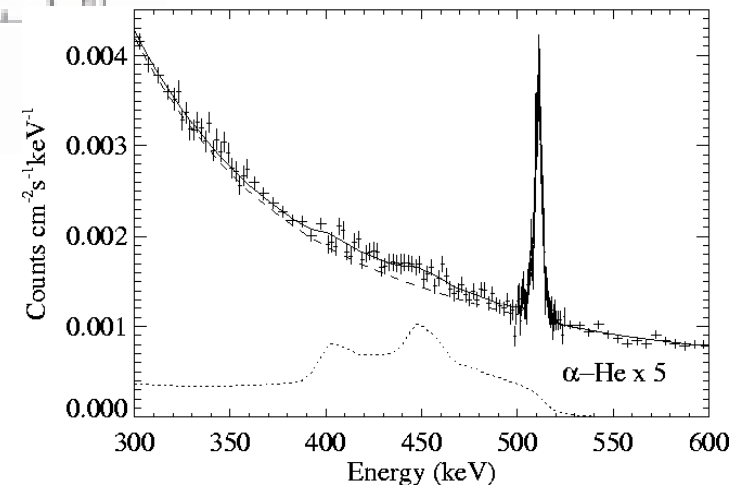
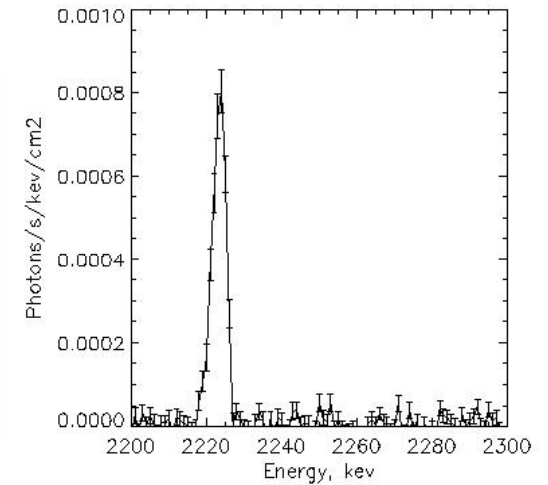


Ion acceleration >30 MeV
is correlated with
relativistic electron
acceleration >300 keV

Ion acceleration >30 MeV
is poorly correlated with
electron acceleration >50
keV, with the possibility of
two separate classes of
flares
(Shih et al, 2007)



October 28, 2003 X-class flare (Share et al, 2004) spectrum

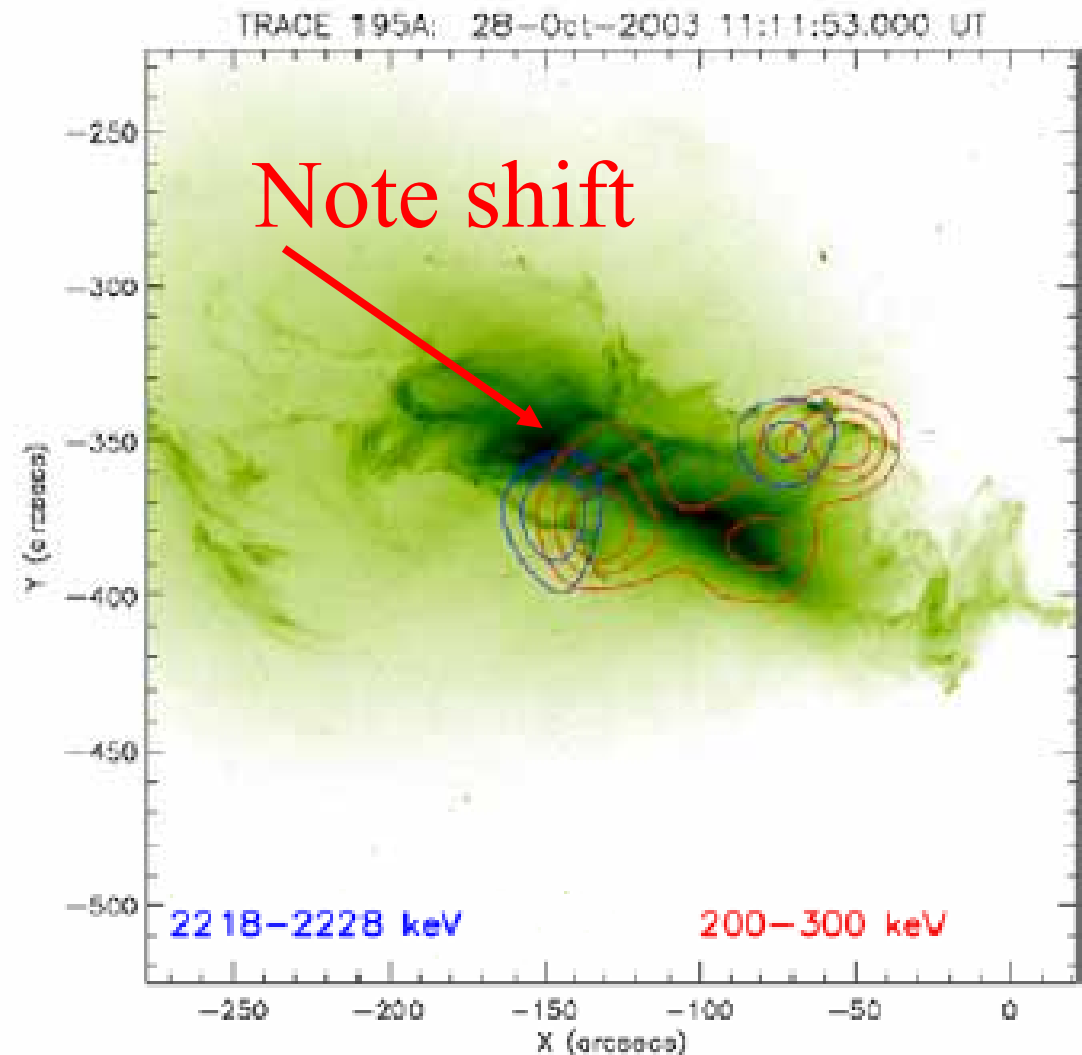


Alpha-alpha lines favours forward **isotropic distribution**

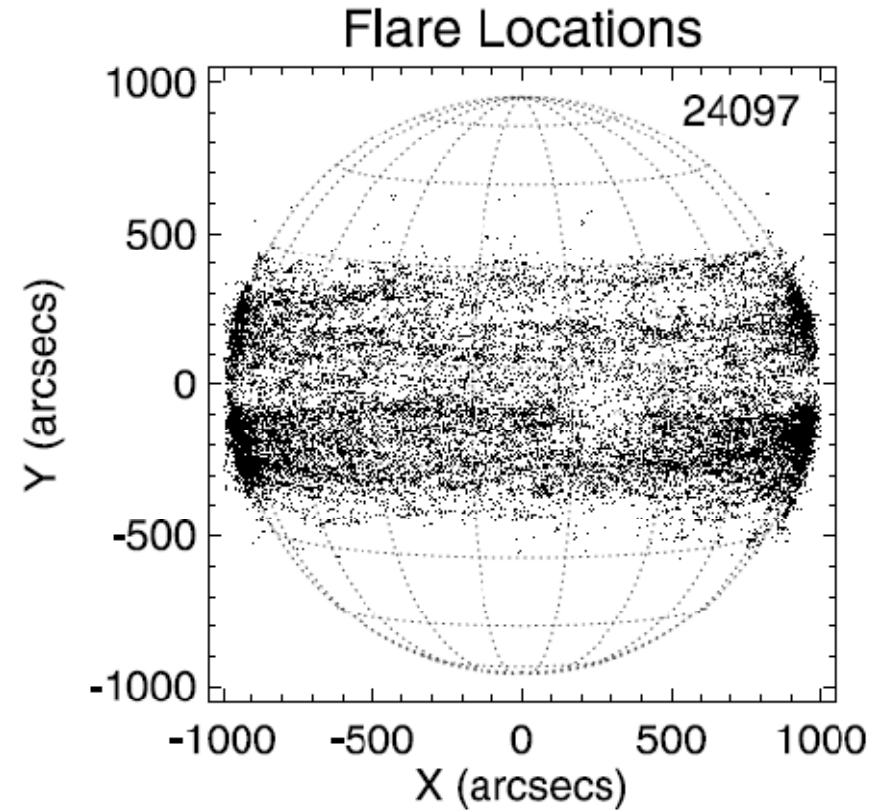
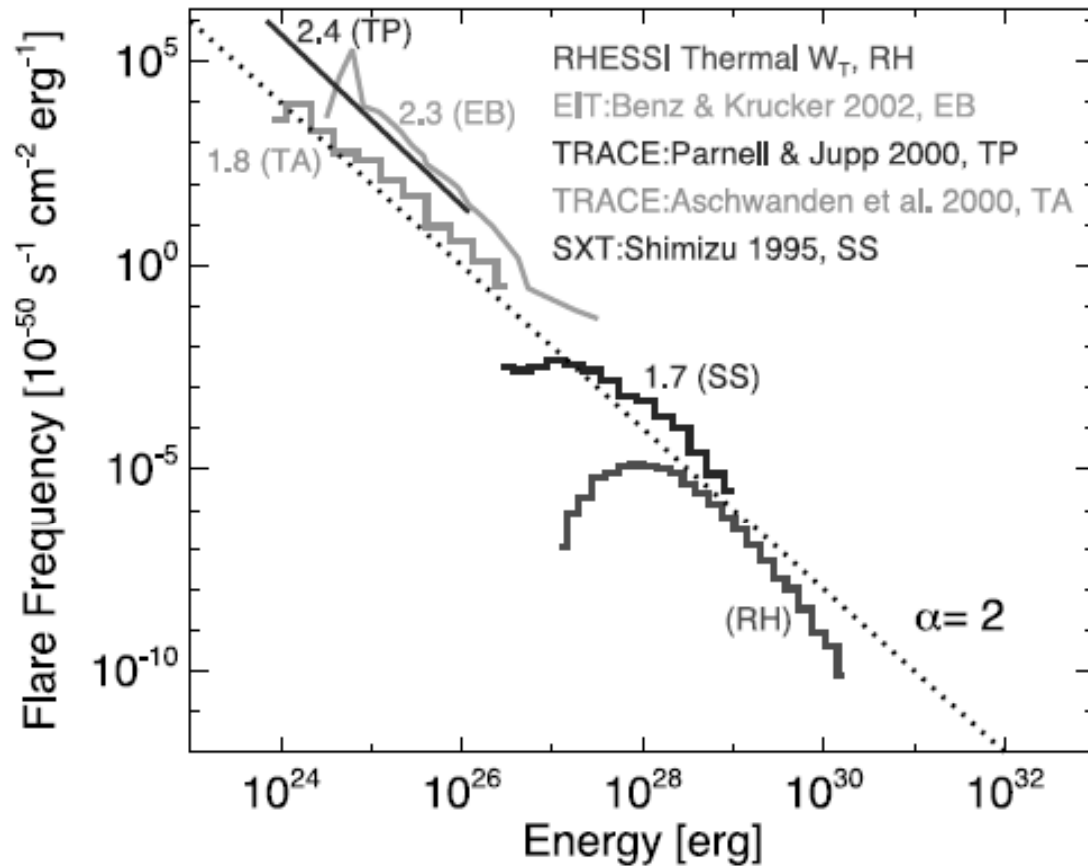
Proton & Alpha power law index is 3.75

2.2MeV line shows ~100 s delay

Hurford et al, 2006: Imaging of the 2.223 MeV neutroncapture line (blue contours) and the HXR electron bremsstrahlung (red contours) of the flare on October 28, 2003. The underlying image is from TRACE at 195 Å. The X-ray and γ -ray imaging shown here used exactly the same selection of detector arrays and imaging procedure. Note the apparent loop-top source in the hard X-ray contours.



Why do electrons and ions emit in different locations?



Christe et al, 2008

Hunnah et al, 2008

Do the observations rule out microflare/nanoflare heating scenario?

RHESSI X-ray spectroscopic data allow to scrutinise current electron acceleration/propagation models.

Spatially resolved electron spectra help to understand the physics of electron transport/acceleration - **Do we understand particle transport?**

Non-thermal hard X-ray emission from coronal sources: **Electron trapping or the signature of acceleration?**

If the electron distribution has a lower value low-energy cutoff (<12 keV), do we systematically **underestimate the total number of accelerated electrons** ?

Anisotropy of electrons: How do the electrons propagate downward but have **close to isotropic electron distribution** ? Propagation effects or electron acceleration is extended ?

2.2MeV line sources are displaced with respect to X-ray sources for 2 events. **Are the electrons and ions are accelerated in different regions?**